

FIFTH EDITION

# Differential Equations

## Computing and Modeling

C. HENRY EDWARDS  
DAVID E. PENNEY  
DAVID T. CALVIS

# DIFFERENTIAL EQUATIONS

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**Computing and Modeling**

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Fifth Edition

**C. Henry Edwards**

**David E. Penney**

*The University of Georgia*

*with the assistance of*

**David Calvis**

*Baldwin Wallace College*

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# APPLICATION MODULES

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The modules listed here follow the indicated sections in the text. Most provide computing projects that illustrate the content of the corresponding text sections. *Maple*, *Mathematica*, and *MATLAB* versions of these investigations are included in the Applications Manual that accompanies this text.

- 1.3 Computer-Generated Slope Fields and Solution Curves
- 1.4 The Logistic Equation
- 1.5 Indoor Temperature Oscillations
- 1.6 Computer Algebra Solutions
- 2.1 Logistic Modeling of Population Data
- 2.3 Rocket Propulsion
- 2.4 Implementing Euler's Method
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- 3.1 Plotting Second-Order Solution Families
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- 3.6 Forced Vibrations
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- 10.1 Numerical Eigenfunction Expansions
- 10.2 Numerical Heat Flow Investigations
- 10.3 Vibrating Beams and Diving Boards
- 10.4 Bessel Functions and Heated Cylinders



# P R E F A C E

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This is a textbook for the standard introductory differential equations course taken by science and engineering students. Its updated content reflects the wide availability of technical computing environments like *Maple*, *Mathematica*, and MATLAB that now are used extensively by practicing engineers and scientists. The traditional manual and symbolic methods are augmented with coverage of qualitative and computer-based methods that employ numerical computation and graphical visualization to develop greater conceptual understanding. A bonus of this more comprehensive approach is accessibility to a wider range of more realistic applications of differential equations.

## Principal Features of This Revision

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This 5th edition is a comprehensive and wide-ranging revision.

In addition to fine-tuning the exposition (both text and graphics) in numerous sections throughout the book, new applications have been inserted (including biological), and we have exploited throughout the new interactive computer technology that is now available to students on devices ranging from desktop and laptop computers to smart phones and graphing calculators. It also utilizes computer algebra systems such as *Mathematica*, Maple, and MATLAB as well as online web sites such as Wolfram|Alpha.

However, with a single exception of a new section inserted in Chapter 5 (noted below), the classtested table of contents of the book remains unchanged. Therefore, instructors' notes and syllabi will not require revision to continue teaching with this new edition.

A conspicuous feature of this edition is the insertion of about 75 new computer-generated figures, many of them illustrating how interactive computer applications with slider bars or touchpad controls can be used to change initial values or parameters in a differential equation, allowing the user to immediately see in real time the resulting changes in the structure of its solutions.

Some illustrations of the various types of revision and updating exhibited in this edition:

**New Interactive Technology and Graphics** New figures inserted throughout illustrate the facility offered by modern computing technology platforms for the user to interactively vary initial conditions and other parameters in real time. Thus, using a mouse or touchpad, the initial point for an initial value problem can be dragged to a new location, and the corresponding solution curve is automatically redrawn and dragged along with its initial point. For instance, see the application modules for Sections 1.3 (page 28) and 3.1 (page 148). Using slider bars in an interactive graphic, the coefficients or other parameters in a linear system can be varied, and the corresponding changes in its direction field and phase plane portrait are automatically shown; see the Section 5.3 application module (page 319).

**New Exposition** In a number of sections, new text and graphics have been inserted to enhance student understanding of the subject matter. For instance, see the new introductory treatments of separable equations in Section 1.4 (page 30), of linear equations in Section 1.5 (page 45), of isolated critical points in Sections 6.1 (page 372) and 6.2 (page 383) and the new example in Section 9.6 (on page 618) showing a vibrating string with a momentary “flat spot”. Examples and accompanying graphics have been updated in Sections 2.4-2.6 and 4.2-4.3 to illustrate new graphing calculators.

**New Content** The single entirely new section for this edition is Section 5.3, which is devoted to the construction of a “gallery” of phase plane portraits illustrating all the possible geometric behaviors of solutions of the 2-dimensional linear system  $\mathbf{x} = \mathbf{Ax}$ . In motivation and preparation for the detailed study of eigenvalue-eigenvector methods in subsequent sections of Chapter 5 (which then follow in the same order as in the previous edition), Section 5.3 shows how the particular arrangements of eigenvalues and eigenvectors of the coefficient matrix  $\mathbf{A}$  correspond to identifiable patterns—“fingerprints,” so to speak—in the phase plane portrait of the system  $\mathbf{x} = \mathbf{Ax}$ . The resulting gallery is shown in the two pages of phase plane portraits that comprise Figure 5.3.16 (pages 315–316) at the end of the section. The new 5.3 application module (on dynamic phase plane portraits, page 319) shows how students can use interactive computer systems to “bring to life” this gallery, by allowing initial conditions, eigenvalues, and even eigenvectors to vary in real time. This dynamic approach is then illustrated with several new graphics inserted in the remainder of Chapter 5. Finally, for a new biological application, see the application module for Section 6.4, which now includes a substantial investigation (page 423) of the nonlinear FitzHugh-Nagumo equations in neuroscience, which were introduced to model the behavior of neurons in the nervous system.

## Computing Features

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The following features highlight the computing technology that distinguishes much of our exposition.

- Over 500 *computer-generated figures* show students vivid pictures of direction fields, solution curves, and phase plane portraits that bring symbolic solutions of differential equations to life.
- About 30 *application modules* follow key sections throughout the text. Most of these applications outline “technology neutral” investigations illustrating the use of technical computing systems and seek to actively engage students in the application of new technology.
- A fresh *numerical emphasis* that is afforded by the early introduction of numerical solution techniques in Chapter 2 (on mathematical models and numerical methods). Here and in Chapter 4, where numerical techniques for systems are treated, a concrete and tangible flavor is achieved by the inclusion of numerical algorithms presented in parallel fashion for systems ranging from graphing calculators to MATLAB.

## Modeling Features

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Mathematical modeling is a goal and constant motivation for the study of differential equations. To sample the range of applications in this text, take a look at the following questions:

- What explains the commonly observed time lag between indoor and outdoor daily temperature oscillations? (Section 1.5)

- What makes the difference between doomsday and extinction in alligator populations? (Section 2.1)
- How do a unicycle and a twoaxle car react differently to road bumps? (Sections 3.7 and 5.4)
- How can you predict the time of next perihelion passage of a newly observed comet? (Section 4.3)
- Why might an earthquake demolish one building and leave standing the one next door? (Section 5.4)
- What determines whether two species will live harmoniously together, or whether competition will result in the extinction of one of them and the survival of the other? (Section 6.3)
- Why and when does non-linearity lead to chaos in biological and mechanical systems? (Section 6.5)
- If a mass on a spring is periodically struck with a hammer, how does the behavior of the mass depend on the frequency of the hammer blows? (Section 7.6)

## Organization and Content

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We have reshaped the usual approach and sequence of topics to accommodate new technology and new perspectives. For instance:

- After a precis of first-order equations in Chapter 1 (though with the coverage of certain traditional symbolic methods streamlined a bit), Chapter 2 offers an early introduction to mathematical modeling, stability and qualitative properties of differential equations, and numerical methods—a combination of topics that frequently are dispersed later in an introductory course. Chapter 3 includes the standard methods of solution of linear differential equations of higher order, particularly those with constant coefficients, and provides an especially wide range of applications involving simple mechanical systems and electrical circuits; the chapter ends with an elementary treatment of endpoint problems and eigenvalues.
- Chapters 4 and 5 provide a flexible treatment of linear systems. Motivated by current trends in science and engineering education and practice, Chapter 4 offers an early, intuitive introduction to first-order systems, models, and numerical approximation techniques. Chapter 5 begins with a self-contained treatment of the linear algebra that is needed, and then presents the eigenvalue approach to linear systems. It includes a wide range of applications (ranging from railway cars to earthquakes) of all the various cases of the eigenvalue method. Section 5.5 includes a fairly extensive treatment of matrix exponentials, which are exploited in Section 5.6 on nonhomogeneous linear systems.
- Chapter 6 on nonlinear systems and phenomena ranges from phase plane analysis to ecological and mechanical systems to a concluding section on chaos and bifurcation in dynamical systems. Section 6.5 presents an elementary introduction to such contemporary topics as period-doubling in biological and mechanical systems, the pitchfork diagram, and the Lorenz strange attractor (all illustrated with vivid computer graphics).
- Laplace transform methods (Chapter 7) follow the material on linear and nonlinear systems, but can be covered at any earlier point (after Chapter 3) the instructor desires.

This book includes enough material appropriately arranged for different courses varying in length from one quarter to two semesters. The longer version *Differential Equations and Boundary Value Problems: Computing and Modeling*

(0-321-79698-5) contains additional chapters on power series methods, Fourier series, separation of variables and partial differential equations).

## Student and Instructor Resources

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The answer section has been expanded considerably to increase its value as a learning aid. It now includes the answers to most odd-numbered problems plus a good many even-numbered ones. The **Instructor's Solutions Manual** (0-321-79701-9) available at [www.pearsonhighered.com/irc](http://www.pearsonhighered.com/irc) provides worked-out solutions for most of the problems in the book, and the **Student Solutions Manual** (0-321-79700-0) contains solutions for most of the odd-numbered problems. Both manuals have been reworked extensively for this edition with improved explanations and more details inserted in the solutions of many problems.

The approximately 30 application modules in the text contain additional problem and project material designed largely to engage students in the exploration and application of computational technology. These investigations are expanded considerably in the **Applications Manual** (0-321-79704-3) that accompanies the text and supplements it with additional and sometimes more challenging investigations. Each section in this manual has parallel subsections **Using Maple**, **Using Mathematica**, and **Using MATLAB** that detail the applicable methods and techniques of each system, and will afford student users an opportunity to compare the merits and styles of different computational systems. These materials—as well as the text of the **Applications Manual** itself—are freely available at the web site [www.pearsonhighered.com/mathstatsresources](http://www.pearsonhighered.com/mathstatsresources).

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# First-Order Differential Equations

## 1.1 Differential Equations and Mathematical Models

The laws of the universe are written in the language of mathematics. Algebra is sufficient to solve many static problems, but the most interesting natural phenomena involve change and are described by equations that relate changing quantities.

Because the derivative  $dx/dt = f'(t)$  of the function  $f$  is the rate at which the quantity  $x = f(t)$  is changing with respect to the independent variable  $t$ , it is natural that equations involving derivatives are frequently used to describe the changing universe. An equation relating an unknown function and one or more of its derivatives is called a **differential equation**.

**Example 1** The differential equation

$$\frac{dx}{dt} = x^2 + t^2$$

involves both the unknown function  $x(t)$  and its first derivative  $x'(t) = dx/dt$ . The differential equation

$$\frac{d^2y}{dx^2} + 3\frac{dy}{dx} + 7y = 0$$

involves the unknown function  $y$  of the independent variable  $x$  and the first two derivatives  $y'$  and  $y''$  of  $y$ . ■

The study of differential equations has three principal goals:

1. To discover the differential equation that describes a specified physical situation.
2. To find—either exactly or approximately—the appropriate solution of that equation.
3. To interpret the solution that is found.

In algebra, we typically seek the unknown *numbers* that satisfy an equation such as  $x^3 + 7x^2 - 11x + 41 = 0$ . By contrast, in solving a differential equation, we

are challenged to find the unknown *functions*  $y = y(x)$  for which an identity such as  $y'(x) = 2xy(x)$ —that is, the differential equation

$$\frac{dy}{dx} = 2xy$$

—holds on some interval of real numbers. Ordinarily, we will want to find *all* solutions of the differential equation, if possible.

**Example 2** If  $C$  is a constant and

$$y(x) = Ce^{x^2}, \quad (1)$$

then

$$\frac{dy}{dx} = C(2xe^{x^2}) = (2x)(Ce^{x^2}) = 2xy.$$

Thus every function  $y(x)$  of the form in Eq. (1) *satisfies*—and thus is a solution of—the differential equation

$$\frac{dy}{dx} = 2xy \quad (2)$$

for all  $x$ . In particular, Eq. (1) defines an *infinite* family of different solutions of this differential equation, one for each choice of the arbitrary constant  $C$ . By the method of separation of variables (Section 1.4) it can be shown that every solution of the differential equation in (2) is of the form in Eq. (1). ■

## Differential Equations and Mathematical Models

The following three examples illustrate the process of translating scientific laws and principles into differential equations. In each of these examples the independent variable is time  $t$ , but we will see numerous examples in which some quantity other than time is the independent variable.

**Example 3** Newton's law of cooling may be stated in this way: The *time rate of change* (the rate of change with respect to time  $t$ ) of the temperature  $T(t)$  of a body is proportional to the difference between  $T$  and the temperature  $A$  of the surrounding medium (Fig. 1.1.1). That is,

$$\frac{dT}{dt} = -k(T - A), \quad (3)$$

where  $k$  is a positive constant. Observe that if  $T > A$ , then  $dT/dt < 0$ , so the temperature is a decreasing function of  $t$  and the body is cooling. But if  $T < A$ , then  $dT/dt > 0$ , so that  $T$  is increasing.

Thus the physical law is translated into a differential equation. If we are given the values of  $k$  and  $A$ , we should be able to find an explicit formula for  $T(t)$ , and then—with the aid of this formula—we can predict the future temperature of the body. ■

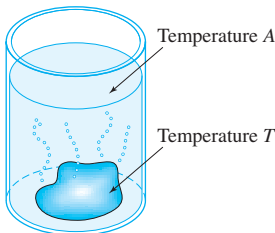
**Example 4** Torricelli's law implies that the *time rate of change* of the volume  $V$  of water in a draining tank (Fig. 1.1.2) is proportional to the square root of the depth  $y$  of water in the tank:

$$\frac{dV}{dt} = -k\sqrt{y}, \quad (4)$$

where  $k$  is a constant. If the tank is a cylinder with vertical sides and cross-sectional area  $A$ , then  $V = Ay$ , so  $dV/dt = A \cdot (dy/dt)$ . In this case Eq. (4) takes the form

$$\frac{dy}{dt} = -h\sqrt{y}, \quad (5)$$

where  $h = k/A$  is a constant. ■



**FIGURE 1.1.1.** Newton's law of cooling, Eq. (3), describes the cooling of a hot rock in water.

**Example 5**

The *time rate of change* of a population  $P(t)$  with constant birth and death rates is, in many simple cases, proportional to the size of the population. That is,

$$\frac{dP}{dt} = kP, \quad (6)$$

where  $k$  is the constant of proportionality. ■

Let us discuss Example 5 further. Note first that each function of the form

$$P(t) = Ce^{kt} \quad (7)$$

is a solution of the differential equation

$$\frac{dP}{dt} = kP$$

in (6). We verify this assertion as follows:

$$P'(t) = Cke^{kt} = k(Ce^{kt}) = kP(t)$$

for all real numbers  $t$ . Because substitution of each function of the form given in (7) into Eq. (6) produces an identity, all such functions are solutions of Eq. (6).

Thus, even if the value of the constant  $k$  is known, the differential equation  $dP/dt = kP$  has *infinitely many* different solutions of the form  $P(t) = Ce^{kt}$ , one for each choice of the “arbitrary” constant  $C$ . This is typical of differential equations. It is also fortunate, because it may allow us to use additional information to select from among all these solutions a particular one that fits the situation under study.

**Example 6**

Suppose that  $P(t) = Ce^{kt}$  is the population of a colony of bacteria at time  $t$ , that the population at time  $t = 0$  (hours, h) was 1000, and that the population doubled after 1 h. This additional information about  $P(t)$  yields the following equations:

$$\begin{aligned} 1000 &= P(0) = Ce^0 = C, \\ 2000 &= P(1) = Ce^k. \end{aligned}$$

It follows that  $C = 1000$  and that  $e^k = 2$ , so  $k = \ln 2 \approx 0.693147$ . With this value of  $k$  the differential equation in (6) is

$$\frac{dP}{dt} = (\ln 2)P \approx (0.693147)P.$$

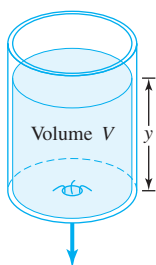
Substitution of  $k = \ln 2$  and  $C = 1000$  in Eq. (7) yields the particular solution

$$P(t) = 1000e^{(\ln 2)t} = 1000(e^{\ln 2})^t = 1000 \cdot 2^t \quad (\text{because } e^{\ln 2} = 2)$$

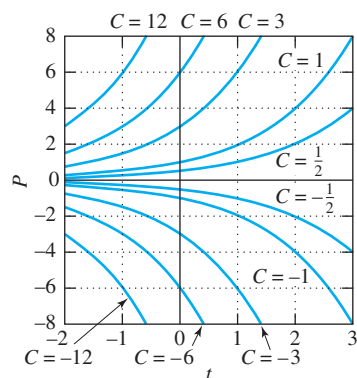
that satisfies the given conditions. We can use this particular solution to predict future populations of the bacteria colony. For instance, the predicted number of bacteria in the population after one and a half hours (when  $t = 1.5$ ) is

$$P(1.5) = 1000 \cdot 2^{3/2} \approx 2828. \quad \blacksquare$$

The condition  $P(0) = 1000$  in Example 6 is called an **initial condition** because we frequently write differential equations for which  $t = 0$  is the “starting time.” Figure 1.1.3 shows several different graphs of the form  $P(t) = Ce^{kt}$  with  $k = \ln 2$ . The graphs of all the infinitely many solutions of  $dP/dt = kP$  in fact fill the entire two-dimensional plane, and no two intersect. Moreover, the selection of any one solution passes through each such point, we see in this case that an initial condition  $P(0) = P_0$  determines a unique solution agreeing with the given data.



**FIGURE 1.1.2.** Newton’s law of cooling, Eq. (3), describes the cooling of a hot rock in water.



**FIGURE 1.1.3.** Graphs of  $P(t) = Ce^{kt}$  with  $k = \ln 2$ .



## Mathematical Models

Our brief discussion of population growth in Examples 5 and 6 illustrates the crucial process of *mathematical modeling* (Fig. 1.1.4), which involves the following:

1. The formulation of a real-world problem in mathematical terms; that is, the construction of a mathematical model.
2. The analysis or solution of the resulting mathematical problem.
3. The interpretation of the mathematical results in the context of the original real-world situation—for example, answering the question originally posed.

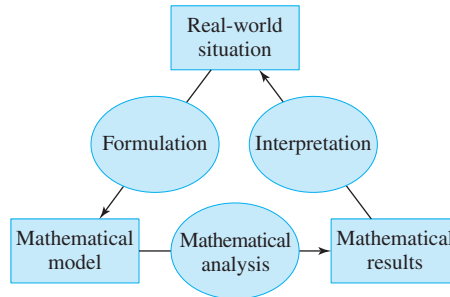


FIGURE 1.1.4. The process of mathematical modeling.

In the population example, the real-world problem is that of determining the population at some future time. A **mathematical model** consists of a list of variables ( $P$  and  $t$ ) that describe the given situation, together with one or more equations relating these variables ( $dP/dt = kP$ ,  $P(0) = P_0$ ) that are known or are assumed to hold. The mathematical analysis consists of solving these equations (here, for  $P$  as a function of  $t$ ). Finally, we apply these mathematical results to attempt to answer the original real-world question.

As an example of this process, think of first formulating the mathematical model consisting of the equations  $dP/dt = kP$ ,  $P(0) = 1000$ , describing the bacteria population of Example 6. Then our mathematical analysis there consisted of solving for the solution function  $P(t) = 1000e^{(\ln 2)t} = 1000 \cdot 2^t$  as our mathematical result. For an interpretation in terms of our real-world situation—the actual bacteria population—we substituted  $t = 1.5$  to obtain the predicted population of  $P(1.5) \approx 2828$  bacteria after 1.5 hours. If, for instance, the bacteria population is growing under ideal conditions of unlimited space and food supply, our prediction may be quite accurate, in which case we conclude that the mathematical model is adequate for studying this particular population.

On the other hand, it may turn out that no solution of the selected differential equation accurately fits the actual population we're studying. For instance, for *no* choice of the constants  $C$  and  $k$  does the solution  $P(t) = Ce^{kt}$  in Eq. (7) accurately describe the actual growth of the human population of the world over the past few centuries. We must conclude that the differential equation  $dP/dt = kP$  is inadequate for modeling the world population—which in recent decades has “leveled off” as compared with the steeply climbing graphs in the upper half ( $P > 0$ ) of Fig. 1.1.3. With sufficient insight, we might formulate a new mathematical model including a perhaps more complicated differential equation, one that takes into account such factors as a limited food supply and the effect of increased population on birth and death rates. With the formulation of this new mathematical model, we may attempt to traverse once again the diagram of Fig. 1.1.4 in a counterclockwise manner. If we can solve the new differential equation, we get new solution functions to com-

pare with the real-world population. Indeed, a successful population analysis may require refining the mathematical model still further as it is repeatedly measured against real-world experience.

But in Example 6 we simply ignored any complicating factors that might affect our bacteria population. This made the mathematical analysis quite simple, perhaps unrealistically so. A satisfactory mathematical model is subject to two contradictory requirements: It must be sufficiently detailed to represent the real-world situation with relative accuracy, yet it must be sufficiently simple to make the mathematical analysis practical. If the model is so detailed that it fully represents the physical situation, then the mathematical analysis may be too difficult to carry out. If the model is too simple, the results may be so inaccurate as to be useless. Thus there is an inevitable tradeoff between what is physically realistic and what is mathematically possible. The construction of a model that adequately bridges this gap between realism and feasibility is therefore the most crucial and delicate step in the process. Ways must be found to simplify the model mathematically without sacrificing essential features of the real-world situation.

Mathematical models are discussed throughout this book. The remainder of this introductory section is devoted to simple examples and to standard terminology used in discussing differential equations and their solutions.

### Examples and Terminology

**Example 7** If  $C$  is a constant and  $y(x) = 1/(C - x)$ , then

$$\frac{dy}{dx} = \frac{1}{(C - x)^2} = y^2$$

if  $x \neq C$ . Thus

$$y(x) = \frac{1}{C - x} \quad (8)$$

defines a solution of the differential equation

$$\frac{dy}{dx} = y^2 \quad (9)$$

on any interval of real numbers not containing the point  $x = C$ . Actually, Eq. (8) defines a *one-parameter family* of solutions of  $dy/dx = y^2$ , one for each value of the arbitrary constant or “parameter”  $C$ . With  $C = 1$  we get the particular solution

$$y(x) = \frac{1}{1 - x}$$

that satisfies the initial condition  $y(0) = 1$ . As indicated in Fig. 1.1.5, this solution is continuous on the interval  $(-\infty, 1)$  but has a vertical asymptote at  $x = 1$ . ■

**Example 8** Verify that the function  $y(x) = 2x^{1/2} - x^{1/2} \ln x$  satisfies the differential equation

$$4x^2 y'' + y = 0 \quad (10)$$

for all  $x > 0$ .

**Solution** First we compute the derivatives

$$y'(x) = -\frac{1}{2}x^{-1/2} \ln x \quad \text{and} \quad y''(x) = \frac{1}{4}x^{-3/2} \ln x - \frac{1}{2}x^{-3/2}.$$

Then substitution into Eq. (10) yields

$$4x^2 y'' + y = 4x^2 \left( \frac{1}{4}x^{-3/2} \ln x - \frac{1}{2}x^{-3/2} \right) + 2x^{1/2} - x^{1/2} \ln x = 0$$

if  $x$  is positive, so the differential equation is satisfied for all  $x > 0$ . ■

The fact that we can write a differential equation is not enough to guarantee that it has a solution. For example, it is clear that the differential equation

$$(y')^2 + y^2 = -1 \tag{11}$$

has *no* (real-valued) solution, because the sum of nonnegative numbers cannot be negative. For a variation on this theme, note that the equation

$$(y')^2 + y^2 = 0 \tag{12}$$

obviously has only the (real-valued) solution  $y(x) \equiv 0$ . In our previous examples any differential equation having at least one solution indeed had infinitely many.

The **order** of a differential equation is the order of the highest derivative that appears in it. The differential equation of Example 8 is of second order, those in Examples 2 through 7 are first-order equations, and

$$y^{(4)} + x^2y^{(3)} + x^5y = \sin x$$

is a fourth-order equation. The most general form of an ***n*th-order** differential equation with independent variable  $x$  and unknown function or dependent variable  $y = y(x)$  is

$$F(x, y, y', y'', \dots, y^{(n)}) = 0, \tag{13}$$

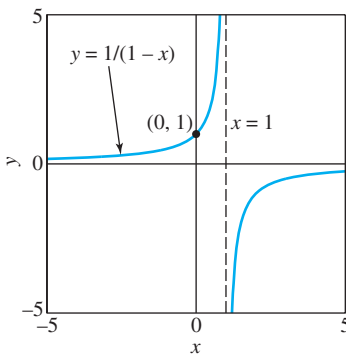
where  $F$  is a specific real-valued function of  $n + 2$  variables.

Our use of the word *solution* has been until now somewhat informal. To be precise, we say that the continuous function  $u = u(x)$  is a **solution** of the differential equation in (13) **on the interval**  $I$  provided that the derivatives  $u', u'', \dots, u^{(n)}$  exist on  $I$  and

$$F(x, u, u', u'', \dots, u^{(n)}) = 0$$

for all  $x$  in  $I$ . For the sake of brevity, we may say that  $u = u(x)$  **satisfies** the differential equation in (13) on  $I$ .

**Remark** Recall from elementary calculus that a differentiable function on an open interval is necessarily continuous there. This is why only a continuous function can qualify as a (differentiable) solution of a differential equation on an interval. ■



**FIGURE 1.1.5.** The solution of  $y' = y^2$  defined by  $y(x) = 1/(1 - x)$ .

**Example 7**  
Continued

Figure 1.1.5 shows the two “connected” branches of the graph  $y = 1/(1 - x)$ . The left-hand branch is the graph of a (continuous) solution of the differential equation  $y' = y^2$  that is defined on the interval  $(-\infty, 1)$ . The right-hand branch is the graph of a *different* solution of the differential equation that is defined (and continuous) on the different interval  $(1, \infty)$ . So the single formula  $y(x) = 1/(1 - x)$  actually defines two different solutions (with different domains of definition) of the same differential equation  $y' = y^2$ . ■

**Example 9**

If  $A$  and  $B$  are constants and

$$y(x) = A \cos 3x + B \sin 3x, \tag{14}$$

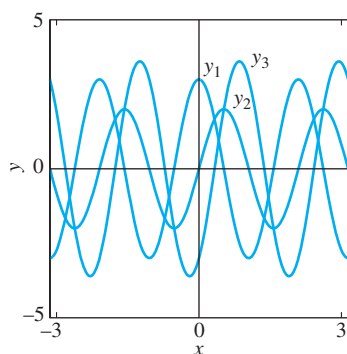
then two successive differentiations yield

$$\begin{aligned} y'(x) &= -3A \sin 3x + 3B \cos 3x, \\ y''(x) &= -9A \cos 3x - 9B \sin 3x = -9y(x) \end{aligned}$$

for all  $x$ . Consequently, Eq. (14) defines what it is natural to call a *two-parameter family* of solutions of the second-order differential equation

$$y'' + 9y = 0 \tag{15}$$

on the whole real number line. Figure 1.1.6 shows the graphs of several such solutions. ■



**FIGURE 1.1.6.** The three solutions  $y_1(x) = 3 \cos 3x$ ,  $y_2(x) = 2 \sin 3x$ , and  $y_3(x) = -3 \cos 3x + 2 \sin 3x$  of the differential equation  $y'' + 9y = 0$ .

Although the differential equations in (11) and (12) are exceptions to the general rule, we will see that an  $n$ th-order differential equation ordinarily has an  $n$ -parameter family of solutions—one involving  $n$  different arbitrary constants or parameters.

In both Eqs. (11) and (12), the appearance of  $y'$  as an implicitly defined function causes complications. For this reason, we will ordinarily assume that any differential equation under study can be solved explicitly for the highest derivative that appears; that is, that the equation can be written in the so-called *normal form*

$$y^{(n)} = G\left(x, y, y', y'', \dots, y^{(n-1)}\right), \quad (16)$$

where  $G$  is a real-valued function of  $n + 1$  variables. In addition, we will always seek only real-valued solutions unless we warn the reader otherwise.

All the differential equations we have mentioned so far are **ordinary** differential equations, meaning that the unknown function (dependent variable) depends on only a *single* independent variable. If the dependent variable is a function of two or more independent variables, then partial derivatives are likely to be involved; if they are, the equation is called a **partial** differential equation. For example, the temperature  $u = u(x, t)$  of a long thin uniform rod at the point  $x$  at time  $t$  satisfies (under appropriate simple conditions) the partial differential equation

$$\frac{\partial u}{\partial t} = k \frac{\partial^2 u}{\partial x^2},$$

where  $k$  is a constant (called the *thermal diffusivity* of the rod). In Chapters 1 through 8 we will be concerned only with *ordinary* differential equations and will refer to them simply as differential equations.

In this chapter we concentrate on *first-order* differential equations of the form

$$\frac{dy}{dx} = f(x, y). \quad (17)$$

We also will sample the wide range of applications of such equations. A typical mathematical model of an applied situation will be an **initial value problem**, consisting of a differential equation of the form in (17) together with an **initial condition**  $y(x_0) = y_0$ . Note that we call  $y(x_0) = y_0$  an initial condition whether or not  $x_0 = 0$ . To **solve** the initial value problem

$$\frac{dy}{dx} = f(x, y), \quad y(x_0) = y_0 \quad (18)$$

means to find a differentiable function  $y = y(x)$  that satisfies both conditions in Eq. (18) on some interval containing  $x_0$ .

### Example 10

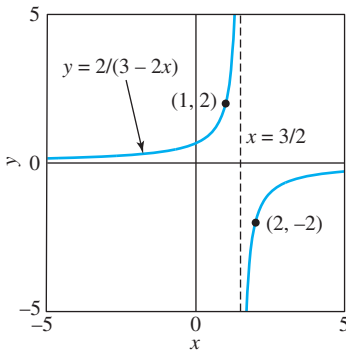
Given the solution  $y(x) = 1/(C - x)$  of the differential equation  $dy/dx = y^2$  discussed in Example 7, solve the initial value problem

$$\frac{dy}{dx} = y^2, \quad y(1) = 2.$$

### Solution

We need only find a value of  $C$  so that the solution  $y(x) = 1/(C - x)$  satisfies the initial condition  $y(1) = 2$ . Substitution of the values  $x = 1$  and  $y = 2$  in the given solution yields

$$2 = y(1) = \frac{1}{C - 1},$$



**FIGURE 1.1.7.** The solutions of  $y' = y^2$  defined by  $y(x) = 2/(3 - 2x)$ .

so  $2C - 2 = 1$ , and hence  $C = \frac{3}{2}$ . With this value of  $C$  we obtain the desired solution

$$y(x) = \frac{1}{\frac{3}{2} - x} = \frac{2}{3 - 2x}.$$

Figure 1.1.7 shows the two branches of the graph  $y = 2/(3 - 2x)$ . The left-hand branch is the graph on  $(-\infty, \frac{3}{2})$  of the solution of the given initial value problem  $y' = y^2$ ,  $y(1) = 2$ . The right-hand branch passes through the point  $(2, -2)$  and is therefore the graph on  $(\frac{3}{2}, \infty)$  of the solution of the different initial value problem  $y' = y^2$ ,  $y(2) = -2$ . ■

The central question of greatest immediate interest to us is this: If we are given a differential equation known to have a solution satisfying a given initial condition, how do we actually *find* or *compute* that solution? And, once found, what can we do with it? We will see that a relatively few simple techniques—separation of variables (Section 1.4), solution of linear equations (Section 1.5), elementary substitution methods (Section 1.6)—are enough to enable us to solve a variety of first-order equations having impressive applications.

## 1.1 Problems

In Problems 1 through 12, verify by substitution that each given function is a solution of the given differential equation. Throughout these problems, primes denote derivatives with respect to  $x$ .

- $y' = 3x^2$ ;  $y = x^3 + 7$
- $y' + 2y = 0$ ;  $y = 3e^{-2x}$
- $y'' + 4y = 0$ ;  $y_1 = \cos 2x$ ,  $y_2 = \sin 2x$
- $y'' = 9y$ ;  $y_1 = e^{3x}$ ,  $y_2 = e^{-3x}$
- $y' = y + 2e^{-x}$ ;  $y = e^x - e^{-x}$
- $y'' + 4y' + 4y = 0$ ;  $y_1 = e^{-2x}$ ,  $y_2 = xe^{-2x}$
- $y'' - 2y' + 2y = 0$ ;  $y_1 = e^x \cos x$ ,  $y_2 = e^x \sin x$
- $y'' + y = 3 \cos 2x$ ,  $y_1 = \cos x - \cos 2x$ ,  $y_2 = \sin x - \cos 2x$
- $y' + 2xy^2 = 0$ ;  $y = \frac{1}{1 + x^2}$
- $x^2 y'' + xy' - y = \ln x$ ;  $y_1 = x - \ln x$ ,  $y_2 = \frac{1}{x} - \ln x$
- $x^2 y'' + 5xy' + 4y = 0$ ;  $y_1 = \frac{1}{x^2}$ ,  $y_2 = \frac{\ln x}{x^2}$
- $x^2 y'' - xy' + 2y = 0$ ;  $y_1 = x \cos(\ln x)$ ,  $y_2 = x \sin(\ln x)$

In Problems 13 through 16, substitute  $y = e^{rx}$  into the given differential equation to determine all values of the constant  $r$  for which  $y = e^{rx}$  is a solution of the equation.

- $3y' = 2y$
- $4y'' = y$
- $y'' + y' - 2y = 0$
- $3y'' + 3y' - 4y = 0$

In Problems 17 through 26, first verify that  $y(x)$  satisfies the given differential equation. Then determine a value of the constant  $C$  so that  $y(x)$  satisfies the given initial condition. Use a computer or graphing calculator (if desired) to sketch several typical solutions of the given differential equation, and highlight the one that satisfies the given initial condition.

- $y' + y = 0$ ;  $y(x) = Ce^{-x}$ ,  $y(0) = 2$
- $y' = 2y$ ;  $y(x) = Ce^{2x}$ ,  $y(0) = 3$
- $y' = y + 1$ ;  $y(x) = Ce^x - 1$ ,  $y(0) = 5$

- $y' = x - y$ ;  $y(x) = Ce^{-x} + x - 1$ ,  $y(0) = 10$
- $y' + 3x^2 y = 0$ ;  $y(x) = Ce^{-x^3}$ ,  $y(0) = 7$
- $e^y y' = 1$ ;  $y(x) = \ln(x + C)$ ,  $y(0) = 0$
- $x \frac{dy}{dx} + 3y = 2x^5$ ;  $y(x) = \frac{1}{4}x^5 + Cx^{-3}$ ,  $y(2) = 1$
- $xy' - 3y = x^3$ ;  $y(x) = x^3(C + \ln x)$ ,  $y(1) = 17$
- $y' = 3x^2(y^2 + 1)$ ;  $y(x) = \tan(x^3 + C)$ ,  $y(0) = 1$
- $y' + y \tan x = \cos x$ ;  $y(x) = (x + C) \cos x$ ,  $y(\pi) = 0$

In Problems 27 through 31, a function  $y = g(x)$  is described by some geometric property of its graph. Write a differential equation of the form  $dy/dx = f(x, y)$  having the function  $g$  as its solution (or as one of its solutions).

- The slope of the graph of  $g$  at the point  $(x, y)$  is the sum of  $x$  and  $y$ .
- The line tangent to the graph of  $g$  at the point  $(x, y)$  intersects the  $x$ -axis at the point  $(x/2, 0)$ .
- Every straight line normal to the graph of  $g$  passes through the point  $(0, 1)$ . Can you guess what the graph of such a function  $g$  might look like?
- The graph of  $g$  is normal to every curve of the form  $y = x^2 + k$  ( $k$  is a constant) where they meet.
- The line tangent to the graph of  $g$  at  $(x, y)$  passes through the point  $(-y, x)$ .

In Problems 32 through 36, write—in the manner of Eqs. (3) through (6) of this section—a differential equation that is a mathematical model of the situation described.

- The time rate of change of a population  $P$  is proportional to the square root of  $P$ .
- The time rate of change of the velocity  $v$  of a coasting motorboat is proportional to the square of  $v$ .
- The acceleration  $dv/dt$  of a Lamborghini is proportional to the difference between 250 km/h and the velocity of the car.



## 1.2 Integrals as General and Particular Solutions

The first-order equation  $dy/dx = f(x, y)$  takes an especially simple form if the right-hand-side function  $f$  does not actually involve the dependent variable  $y$ , so

$$\frac{dy}{dx} = f(x). \quad (1)$$

In this special case we need only integrate both sides of Eq. (1) to obtain

$$y(x) = \int f(x) dx + C. \quad (2)$$

This is a **general solution** of Eq. (1), meaning that it involves an arbitrary constant  $C$ , and for every choice of  $C$  it is a solution of the differential equation in (1). If  $G(x)$  is a particular antiderivative of  $f$ —that is, if  $G'(x) \equiv f(x)$ —then

$$y(x) = G(x) + C. \quad (3)$$

The graphs of any two such solutions  $y_1(x) = G(x) + C_1$  and  $y_2(x) = G(x) + C_2$  on the same interval  $I$  are “parallel” in the sense illustrated by Figs. 1.2.1 and 1.2.2. There we see that the constant  $C$  is geometrically the vertical distance between the two curves  $y(x) = G(x)$  and  $y(x) = G(x) + C$ .

To satisfy an initial condition  $y(x_0) = y_0$ , we need only substitute  $x = x_0$  and  $y = y_0$  into Eq. (3) to obtain  $y_0 = G(x_0) + C$ , so that  $C = y_0 - G(x_0)$ . With this choice of  $C$ , we obtain the **particular solution** of Eq. (1) satisfying the initial value problem

$$\frac{dy}{dx} = f(x), \quad y(x_0) = y_0.$$

We will see that this is the typical pattern for solutions of first-order differential equations. Ordinarily, we will first find a *general solution* involving an arbitrary constant  $C$ . We can then attempt to obtain, by appropriate choice of  $C$ , a *particular solution* satisfying a given initial condition  $y(x_0) = y_0$ .

**Remark** As the term is used in the previous paragraph, a *general solution* of a first-order differential equation is simply a one-parameter family of solutions. A natural question is whether a given general solution contains *every* particular solution of the differential equation. When this is known to be true, we call it **the** general solution of the differential equation. For example, because any two antiderivatives of the same function  $f(x)$  can differ only by a constant, it follows that every solution of Eq. (1) is of the form in (2). Thus Eq. (2) serves to define **the** general solution of (1). ■

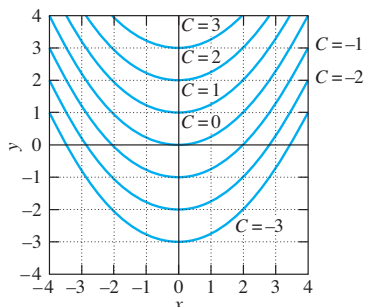


FIGURE 1.2.1. Graphs of  $y = \frac{1}{4}x^2 + C$  for various values of  $C$ .

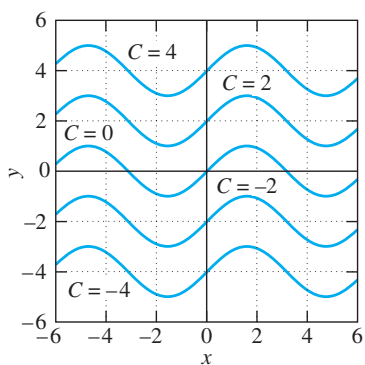


FIGURE 1.2.2. Graphs of  $y = \sin x + C$  for various values of  $C$ .

### Example 1 Solve the initial value problem

$$\frac{dy}{dx} = 2x + 3, \quad y(1) = 2.$$

**Solution** Integration of both sides of the differential equation as in Eq. (2) immediately yields the general solution

$$y(x) = \int (2x + 3) dx = x^2 + 3x + C.$$

Figure 1.2.3 shows the graph  $y = x^2 + 3x + C$  for various values of  $C$ . The particular solution we seek corresponds to the curve that passes through the point  $(1, 2)$ , thereby satisfying the initial condition

$$y(1) = (1)^2 + 3 \cdot (1) + C = 2.$$

It follows that  $C = -2$ , so the desired particular solution is

$$y(x) = x^2 + 3x - 2. \quad \blacksquare$$

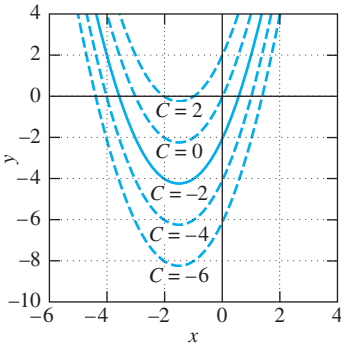


FIGURE 1.2.3. Solution curves for the differential equation in Example 1.

**Second-order equations.** The observation that the special first-order equation  $dy/dx = f(x)$  is readily solvable (provided that an antiderivative of  $f$  can be found) extends to second-order differential equations of the special form

$$\frac{d^2y}{dx^2} = g(x), \quad (4)$$

in which the function  $g$  on the right-hand side involves neither the dependent variable  $y$  nor its derivative  $dy/dx$ . We simply integrate once to obtain

$$\frac{dy}{dx} = \int y''(x) dx = \int g(x) dx = G(x) + C_1,$$

where  $G$  is an antiderivative of  $g$  and  $C_1$  is an arbitrary constant. Then another integration yields

$$y(x) = \int y'(x) dx = \int [G(x) + C_1] dx = \int G(x) dx + C_1x + C_2,$$

where  $C_2$  is a second arbitrary constant. In effect, the second-order differential equation in (4) is one that can be solved by solving successively the *first-order* equations

$$\frac{dv}{dx} = g(x) \quad \text{and} \quad \frac{dy}{dx} = v(x).$$

## Velocity and Acceleration

Direct integration is sufficient to allow us to solve a number of important problems concerning the motion of a particle (or *mass point*) in terms of the forces acting on it. The motion of a particle along a straight line (the  $x$ -axis) is described by its **position function**

$$x = f(t) \quad (5)$$

giving its  $x$ -coordinate at time  $t$ . The **velocity** of the particle is defined to be

$$\blacktriangleright \quad v(t) = f'(t); \quad \text{that is,} \quad v = \frac{dx}{dt}. \quad (6)$$

Its **acceleration**  $a(t)$  is  $a(t) = v'(t) = x''(t)$ ; in Leibniz notation,

$$\blacktriangleright \quad a = \frac{dv}{dt} = \frac{d^2x}{dt^2}. \quad (7)$$

Equation (6) is sometimes applied either in the indefinite integral form  $x(t) = \int v(t) dt$  or in the definite integral form

$$x(t) = x(t_0) + \int_{t_0}^t v(s) ds,$$

which you should recognize as a statement of the fundamental theorem of calculus (precisely because  $dx/dt = v$ ).

Newton's *second law of motion* says that if a force  $F(t)$  acts on the particle and is directed along its line of motion, then

$$ma(t) = F(t); \quad \text{that is,} \quad F = ma, \quad (8)$$



where  $m$  is the mass of the particle. If the force  $F$  is known, then the equation  $x''(t) = F(t)/m$  can be integrated twice to find the position function  $x(t)$  in terms of two constants of integration. These two arbitrary constants are frequently determined by the **initial position**  $x_0 = x(0)$  and the **initial velocity**  $v_0 = v(0)$  of the particle.

**Constant acceleration.** For instance, suppose that the force  $F$ , and therefore the acceleration  $a = F/m$ , are *constant*. Then we begin with the equation

$$\frac{dv}{dt} = a \quad (a \text{ is a constant}) \quad (9)$$

and integrate both sides to obtain

$$v(t) = \int a \, dt = at + C_1.$$

We know that  $v = v_0$  when  $t = 0$ , and substitution of this information into the preceding equation yields the fact that  $C_1 = v_0$ . So

$$v(t) = \frac{dx}{dt} = at + v_0. \quad (10)$$

A second integration gives

$$x(t) = \int v(t) \, dt = \int (at + v_0) \, dt = \frac{1}{2}at^2 + v_0t + C_2,$$

and the substitution  $t = 0$ ,  $x = x_0$  gives  $C_2 = x_0$ . Therefore,

$$x(t) = \frac{1}{2}at^2 + v_0t + x_0. \quad (11)$$

Thus, with Eq. (10) we can find the velocity, and with Eq. (11) the position, of the particle at any time  $t$  in terms of its *constant* acceleration  $a$ , its initial velocity  $v_0$ , and its initial position  $x_0$ .

### Example 2

A lunar lander is falling freely toward the surface of the moon at a speed of 450 meters per second (m/s). Its retrorockets, when fired, provide a constant deceleration of 2.5 meters per second per second ( $\text{m/s}^2$ ) (the gravitational acceleration produced by the moon is assumed to be included in the given deceleration). At what height above the lunar surface should the retrorockets be activated to ensure a “soft touchdown” ( $v = 0$  at impact)?

### Solution

We denote by  $x(t)$  the height of the lunar lander above the surface, as indicated in Fig. 1.2.4. We let  $t = 0$  denote the time at which the retrorockets should be fired. Then  $v_0 = -450$  (m/s, negative because the height  $x(t)$  is decreasing), and  $a = +2.5$ , because an upward thrust increases the velocity  $v$  (although it decreases the *speed*  $|v|$ ). Then Eqs. (10) and (11) become

$$v(t) = 2.5t - 450 \quad (12)$$

and

$$x(t) = 1.25t^2 - 450t + x_0, \quad (13)$$

where  $x_0$  is the height of the lander above the lunar surface at the time  $t = 0$  when the retrorockets should be activated.

From Eq. (12) we see that  $v = 0$  (soft touchdown) occurs when  $t = 450/2.5 = 180$  s (that is, 3 minutes); then substitution of  $t = 180$ ,  $x = 0$  into Eq. (13) yields

$$x_0 = 0 - (1.25)(180)^2 + 450(180) = 40,500$$

meters—that is,  $x_0 = 40.5 \text{ km} \approx 25\frac{1}{6}$  miles. Thus the retrorockets should be activated when the lunar lander is 40.5 kilometers above the surface of the moon, and it will touch down softly on the lunar surface after 3 minutes of decelerating descent. ■

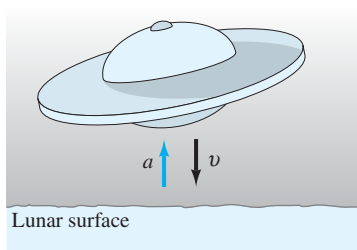


FIGURE 1.2.4. The lunar lander of Example 2.

## Physical Units

Numerical work requires units for the measurement of physical quantities such as distance and time. We sometimes use ad hoc units—such as distance in miles or kilometers and time in hours—in special situations (such as in a problem involving an auto trip). However, the foot-pound-second (fps) and meter-kilogram-second (mks) unit systems are used more generally in scientific and engineering problems. In fact, fps units are commonly used only in the United States (and a few other countries), while mks units constitute the standard international system of scientific units.

	fps units	mks units
Force	pound (lb)	newton (N)
Mass	slug	kilogram (kg)
Distance	foot (ft)	meter (m)
Time	second (s)	second (s)
$g$	$32 \text{ ft/s}^2$	$9.8 \text{ m/s}^2$

The last line of this table gives values for the gravitational acceleration  $g$  at the surface of the earth. Although these approximate values will suffice for most examples and problems, more precise values are  $9.7805 \text{ m/s}^2$  and  $32.088 \text{ ft/s}^2$  (at sea level at the equator).

Both systems are compatible with Newton's second law  $F = ma$ . Thus 1 N is (by definition) the force required to impart an acceleration of  $1 \text{ m/s}^2$  to a mass of 1 kg. Similarly, 1 slug is (by definition) the mass that experiences an acceleration of  $1 \text{ ft/s}^2$  under a force of 1 lb. (We will use mks units in all problems requiring mass units and thus will rarely need slugs to measure mass.)

Inches and centimeters (as well as miles and kilometers) also are commonly used in describing distances. For conversions between fps and mks units it helps to remember that

$$1 \text{ in.} = 2.54 \text{ cm (exactly)} \quad \text{and} \quad 1 \text{ lb} \approx 4.448 \text{ N.}$$

For instance,

$$1 \text{ ft} = 12 \text{ in.} \times 2.54 \frac{\text{cm}}{\text{in.}} = 30.48 \text{ cm,}$$

and it follows that

$$1 \text{ mi} = 5280 \text{ ft} \times 30.48 \frac{\text{cm}}{\text{ft}} = 160934.4 \text{ cm} \approx 1.609 \text{ km.}$$

Thus a posted U.S. speed limit of 50 mi/h means that—in international terms—the legal speed limit is about  $50 \times 1.609 \approx 80.45 \text{ km/h}$ .

## Vertical Motion with Gravitational Acceleration

The **weight**  $W$  of a body is the force exerted on the body by gravity. Substitution of  $a = g$  and  $F = W$  in Newton's second law  $F = ma$  gives

$$W = mg \tag{14}$$

for the weight  $W$  of the mass  $m$  at the surface of the earth (where  $g \approx 32 \text{ ft/s}^2 \approx 9.8 \text{ m/s}^2$ ). For instance, a mass of  $m = 20 \text{ kg}$  has a weight of  $W = (20 \text{ kg})(9.8 \text{ m/s}^2) = 196 \text{ N}$ . Similarly, a mass  $m$  weighing 100 pounds has mks weight

$$W = (100 \text{ lb})(4.448 \text{ N/lb}) = 444.8 \text{ N},$$

so its mass is

$$m = \frac{W}{g} = \frac{444.8 \text{ N}}{9.8 \text{ m/s}^2} \approx 45.4 \text{ kg}.$$

To discuss vertical motion it is natural to choose the  $y$ -axis as the coordinate system for position, frequently with  $y = 0$  corresponding to “ground level.” If we choose the *upward* direction as the positive direction, then the effect of gravity on a vertically moving body is to decrease its height and also to decrease its velocity  $v = dy/dt$ . Consequently, if we ignore air resistance, then the acceleration  $a = dv/dt$  of the body is given by

► 
$$\frac{dv}{dt} = -g. \tag{15}$$

This acceleration equation provides a starting point in many problems involving vertical motion. Successive integrations (as in Eqs. (10) and (11)) yield the velocity and height formulas

$$v(t) = -gt + v_0 \tag{16}$$

and

$$y(t) = -\frac{1}{2}gt^2 + v_0t + y_0. \tag{17}$$

Here,  $y_0$  denotes the initial ( $t = 0$ ) height of the body and  $v_0$  its initial velocity.

**Example 3** (a) Suppose that a ball is thrown straight upward from the ground ( $y_0 = 0$ ) with initial velocity  $v_0 = 96 \text{ (ft/s)}$ , so we use  $g = 32 \text{ ft/s}^2$  in fps units). Then it reaches its maximum height when its velocity (Eq. (16)) is zero,

$$v(t) = -32t + 96 = 0,$$

and thus when  $t = 3 \text{ s}$ . Hence the maximum height that the ball attains is

$$y(3) = -\frac{1}{2} \cdot 32 \cdot 3^2 + 96 \cdot 3 + 0 = 144 \text{ (ft)}$$

(with the aid of Eq. (17)).

(b) If an arrow is shot straight upward from the ground with initial velocity  $v_0 = 49 \text{ (m/s)}$ , so we use  $g = 9.8 \text{ m/s}^2$  in mks units), then it returns to the ground when

$$y(t) = -\frac{1}{2} \cdot (9.8)t^2 + 49t = (4.9)t(-t + 10) = 0,$$

and thus after 10 s in the air. ■

### A Swimmer’s Problem

Figure 1.2.5 shows a northward-flowing river of width  $w = 2a$ . The lines  $x = \pm a$  represent the banks of the river and the  $y$ -axis its center. Suppose that the velocity  $v_R$  at which the water flows increases as one approaches the center of the river, and indeed is given in terms of distance  $x$  from the center by

$$v_R = v_0 \left( 1 - \frac{x^2}{a^2} \right). \tag{18}$$

You can use Eq. (18) to verify that the water does flow the fastest at the center, where  $v_R = v_0$ , and that  $v_R = 0$  at each riverbank.

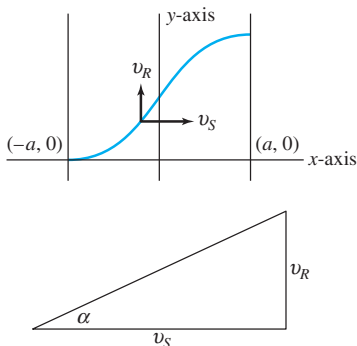


FIGURE 1.2.5. A swimmer’s problem (Example 4).

Suppose that a swimmer starts at the point  $(-a, 0)$  on the west bank and swims due east (relative to the water) with constant speed  $v_S$ . As indicated in Fig. 1.2.5, his velocity vector (relative to the riverbed) has horizontal component  $v_S$  and vertical component  $v_R$ . Hence the swimmer's direction angle  $\alpha$  is given by

$$\tan \alpha = \frac{v_R}{v_S}.$$

Because  $\tan \alpha = dy/dx$ , substitution using (18) gives the differential equation

$$\frac{dy}{dx} = \frac{v_0}{v_S} \left( 1 - \frac{x^2}{a^2} \right) \quad (19)$$

for the swimmer's trajectory  $y = y(x)$  as he crosses the river.

#### Example 4

Suppose that the river is 1 mile wide and that its midstream velocity is  $v_0 = 9$  mi/h. If the swimmer's velocity is  $v_S = 3$  mi/h, then Eq. (19) takes the form

$$\frac{dy}{dx} = 3(1 - 4x^2).$$

Integration yields

$$y(x) = \int (3 - 12x^2) dx = 3x - 4x^3 + C$$

for the swimmer's trajectory. The initial condition  $y(-\frac{1}{2}) = 0$  yields  $C = 1$ , so

$$y(x) = 3x - 4x^3 + 1.$$

Then

$$y\left(\frac{1}{2}\right) = 3\left(\frac{1}{2}\right) - 4\left(\frac{1}{2}\right)^3 + 1 = 2,$$

so the swimmer drifts 2 miles downstream while he swims 1 mile across the river. ■

## 1.2 Problems

In Problems 1 through 10, find a function  $y = f(x)$  satisfying the given differential equation and the prescribed initial condition.

1.  $\frac{dy}{dx} = 2x + 1$ ;  $y(0) = 3$

2.  $\frac{dy}{dx} = (x - 2)^2$ ;  $y(2) = 1$

3.  $\frac{dy}{dx} = \sqrt{x}$ ;  $y(4) = 0$

4.  $\frac{dy}{dx} = \frac{1}{x^2}$ ;  $y(1) = 5$

5.  $\frac{dy}{dx} = \frac{1}{\sqrt{x+2}}$ ;  $y(2) = -1$

6.  $\frac{dy}{dx} = x\sqrt{x^2+9}$ ;  $y(-4) = 0$

7.  $\frac{dy}{dx} = \frac{10}{x^2+1}$ ;  $y(0) = 0$

8.  $\frac{dy}{dx} = \cos 2x$ ;  $y(0) = 1$

9.  $\frac{dy}{dx} = \frac{1}{\sqrt{1-x^2}}$ ;  $y(0) = 0$

10.  $\frac{dy}{dx} = xe^{-x}$ ;  $y(0) = 1$

In Problems 11 through 18, find the position function  $x(t)$  of a moving particle with the given acceleration  $a(t)$ , initial position  $x_0 = x(0)$ , and initial velocity  $v_0 = v(0)$ .

11.  $a(t) = 50$ ,  $v_0 = 10$ ,  $x_0 = 20$

12.  $a(t) = -20$ ,  $v_0 = -15$ ,  $x_0 = 5$

13.  $a(t) = 3t$ ,  $v_0 = 5$ ,  $x_0 = 0$

14.  $a(t) = 2t + 1$ ,  $v_0 = -7$ ,  $x_0 = 4$

15.  $a(t) = 4(t+3)^2$ ,  $v_0 = -1$ ,  $x_0 = 1$

16.  $a(t) = \frac{1}{\sqrt{t+4}}$ ,  $v_0 = -1$ ,  $x_0 = 1$

17.  $a(t) = \frac{1}{(t+1)^3}$ ,  $v_0 = 0$ ,  $x_0 = 0$

18.  $a(t) = 50 \sin 5t$ ,  $v_0 = -10$ ,  $x_0 = 8$

In Problems 19 through 22, a particle starts at the origin and travels along the  $x$ -axis with the velocity function  $v(t)$  whose graph is shown in Figs. 1.2.6 through 1.2.9. Sketch the graph of the resulting position function  $x(t)$  for  $0 \leq t \leq 10$ .

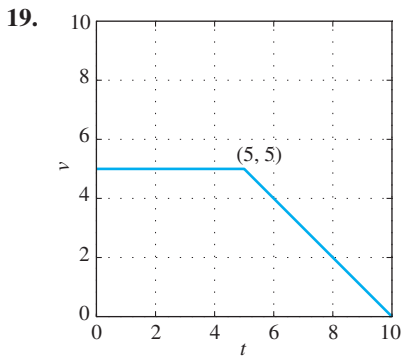


FIGURE 1.2.6. Graph of the velocity function  $v(t)$  of Problem 19.

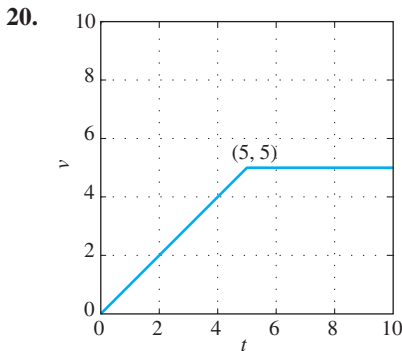


FIGURE 1.2.7. Graph of the velocity function  $v(t)$  of Problem 20.

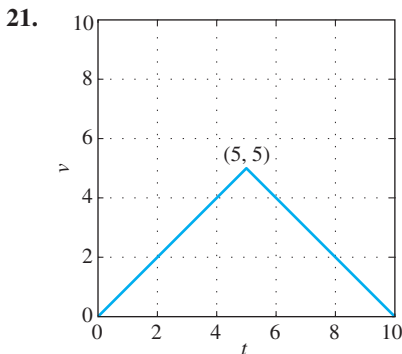


FIGURE 1.2.8. Graph of the velocity function  $v(t)$  of Problem 21.

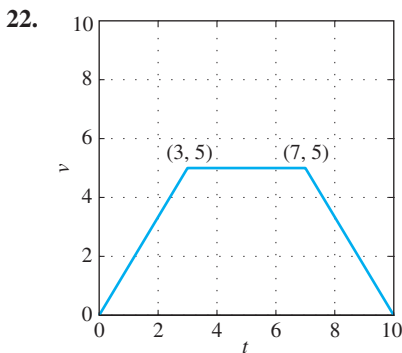


FIGURE 1.2.9. Graph of the velocity function  $v(t)$  of Problem 22.

23. What is the maximum height attained by the arrow of part (b) of Example 3?
24. A ball is dropped from the top of a building 400 ft high. How long does it take to reach the ground? With what speed does the ball strike the ground?
25. The brakes of a car are applied when it is moving at 100 km/h and provide a constant deceleration of 10 meters per second per second ( $m/s^2$ ). How far does the car travel before coming to a stop?
26. A projectile is fired straight upward with an initial velocity of 100 m/s from the top of a building 20 m high and falls to the ground at the base of the building. Find (a) its maximum height above the ground; (b) when it passes the top of the building; (c) its total time in the air.
27. A ball is thrown straight downward from the top of a tall building. The initial speed of the ball is 10 m/s. It strikes the ground with a speed of 60 m/s. How tall is the building?
28. A baseball is thrown straight downward with an initial speed of 40 ft/s from the top of the Washington Monument (555 ft high). How long does it take to reach the ground, and with what speed does the baseball strike the ground?
29. A diesel car gradually speeds up so that for the first 10 s its acceleration is given by

$$\frac{dv}{dt} = (0.12)t^2 + (0.6)t \quad (\text{ft/s}^2).$$

If the car starts from rest ( $x_0 = 0$ ,  $v_0 = 0$ ), find the distance it has traveled at the end of the first 10 s and its velocity at that time.

30. A car traveling at 60 mi/h (88 ft/s) skids 176 ft after its brakes are suddenly applied. Under the assumption that the braking system provides constant deceleration, what is that deceleration? For how long does the skid continue?
31. The skid marks made by an automobile indicated that its brakes were fully applied for a distance of 75 m before it came to a stop. The car in question is known to have a constant deceleration of  $20 \text{ m/s}^2$  under these conditions. How fast—in km/h—was the car traveling when the brakes were first applied?
32. Suppose that a car skids 15 m if it is moving at 50 km/h when the brakes are applied. Assuming that the car has the same constant deceleration, how far will it skid if it is moving at 100 km/h when the brakes are applied?
33. On the planet Gzyx, a ball dropped from a height of 20 ft hits the ground in 2 s. If a ball is dropped from the top of a 200-ft-tall building on Gzyx, how long will it take to hit the ground? With what speed will it hit?
34. A person can throw a ball straight upward from the surface of the earth to a maximum height of 144 ft. How high could this person throw the ball on the planet Gzyx of Problem 33?
35. A stone is dropped from rest at an initial height  $h$  above the surface of the earth. Show that the speed with which it strikes the ground is  $v = \sqrt{2gh}$ .

36. Suppose a woman has enough “spring” in her legs to jump (on earth) from the ground to a height of 2.25 feet. If she jumps straight upward with the same initial velocity on the moon—where the surface gravitational acceleration is (approximately)  $5.3 \text{ ft/s}^2$ —how high above the surface will she rise?
37. At noon a car starts from rest at point  $A$  and proceeds at constant acceleration along a straight road toward point  $B$ . If the car reaches  $B$  at 12:50 P.M. with a velocity of 60 mi/h, what is the distance from  $A$  to  $B$ ?
38. At noon a car starts from rest at point  $A$  and proceeds with constant acceleration along a straight road toward point  $C$ , 35 miles away. If the constantly accelerated car arrives at  $C$  with a velocity of 60 mi/h, at what time does it arrive at  $C$ ?
39. If  $a = 0.5 \text{ mi}$  and  $v_0 = 9 \text{ mi/h}$  as in Example 4, what must the swimmer’s speed  $v_S$  be in order that he drifts only 1 mile downstream as he crosses the river?
40. Suppose that  $a = 0.5 \text{ mi}$ ,  $v_0 = 9 \text{ mi/h}$ , and  $v_S = 3 \text{ mi/h}$  as in Example 4, but that the velocity of the river is given by the fourth-degree function

$$v_R = v_0 \left( 1 - \frac{x^4}{a^4} \right)$$

rather than the quadratic function in Eq. (18). Now find how far downstream the swimmer drifts as he crosses the river.

41. A bomb is dropped from a helicopter hovering at an altitude of 800 feet above the ground. From the ground directly beneath the helicopter, a projectile is fired straight upward toward the bomb, exactly 2 seconds after the bomb is released. With what initial velocity should the projectile be fired in order to hit the bomb at an altitude of exactly 400 feet?
42. A spacecraft is in free fall toward the surface of the moon at a speed of 1000 mph (mi/h). Its retrorockets, when fired, provide a constant deceleration of  $20,000 \text{ mi/h}^2$ . At what height above the lunar surface should the astronauts fire the retrorockets to insure a soft touchdown? (As in Example 2, ignore the moon’s gravitational field.)
43. Arthur Clarke’s *The Wind from the Sun* (1963) describes Diana, a spacecraft propelled by the solar wind. Its aluminized sail provides it with a constant acceleration of  $0.001g = 0.0098 \text{ m/s}^2$ . Suppose this spacecraft starts from rest at time  $t = 0$  and simultaneously fires a projectile (straight ahead in the same direction) that travels at one-tenth of the speed  $c = 3 \times 10^8 \text{ m/s}$  of light. How long will it take the spacecraft to catch up with the projectile, and how far will it have traveled by then?
44. A driver involved in an accident claims he was going only 25 mph. When police tested his car, they found that when its brakes were applied at 25 mph, the car skidded only 45 feet before coming to a stop. But the driver’s skid marks at the accident scene measured 210 feet. Assuming the same (constant) deceleration, determine the speed he was actually traveling just prior to the accident.

## 1.3 Slope Fields and Solution Curves

Consider a differential equation of the form

$$\frac{dy}{dx} = f(x, y) \quad (1)$$

where the right-hand function  $f(x, y)$  involves both the independent variable  $x$  and the dependent variable  $y$ . We might think of integrating both sides in (1) with respect to  $x$ , and hence write  $y(x) = \int f(x, y(x)) dx + C$ . However, this approach does not lead to a solution of the differential equation, because the indicated integral involves the *unknown* function  $y(x)$  itself, and therefore cannot be evaluated explicitly. Actually, there exists *no* straightforward procedure by which a general differential equation can be solved explicitly. Indeed, the solutions of such a simple-looking differential equation as  $y' = x^2 + y^2$  cannot be expressed in terms of the ordinary elementary functions studied in calculus textbooks. Nevertheless, the graphical and numerical methods of this and later sections can be used to construct *approximate* solutions of differential equations that suffice for many practical purposes.

### Slope Fields and Graphical Solutions

There is a simple geometric way to think about solutions of a given differential equation  $y' = f(x, y)$ . At each point  $(x, y)$  of the  $xy$ -plane, the value of  $f(x, y)$  determines a slope  $m = f(x, y)$ . A solution of the differential equation is simply a differentiable function whose graph  $y = y(x)$  has this “correct slope” at each